Asymptotic behavior of energy to Klein-Gordon equations with time dependent potentials

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§1. Scattering for Klein-Gordon equations

We consider the Klein-Gordon equation

$$u_{tt} - \Delta u + m^2 u + V(x, t)u = 0, \quad (x, t) \in \mathbf{R}^n \times \mathbf{R}, \tag{P}$$

where $u_{tt} = \partial^2 u / \partial t^2$, Δ is the *n*-dimensional Laplacian, $n \ge 3$, m = m(x) > 0 and V(x, t) is a potential which is bounded and continuous in $\mathbf{R}^n \times \mathbf{R}$. Now, we define the energy of the solution of this equation at time t:

$$\begin{cases} \|u(t)\|_{E}^{2} = \frac{1}{2} \int_{\mathbf{R}^{n}} \{u_{t}(x,t)^{2} + |\nabla u(x,t)|^{2} + m^{2}u(x,t)^{2}\} dx, \\ \|u(t)\|_{E_{V}}^{2} = \frac{1}{2} \int_{\mathbf{R}^{n}} \{u_{t}(x,t)^{2} + |\nabla u(x,t)|^{2} + m^{2}u(x,t)^{2} + V(x,t)u(x,t)^{2}\} dx. \end{cases}$$

where E is the Hilbert space with the above energy norms.

In this paper, under suitable condition on V(x, t), we shall treat a nondecay of the energy, and develop a scattering theory.

As is easily seen, we have

$$\|u(t)\|_{E_V}^2 - \frac{1}{2} \int_s^t \int_{\mathbf{R}^n} V_\tau(x,\tau) u^2 dx d\tau = \|u(s)\|_{E_V}^2$$

for any t > s. If $V_t(x,t) \leq 0$, then $||u(t)||_{E_V}$ is decreasing with t, and a question rises whether it decays or not as t goes to infinity. We conjecture that the perturbed solutions behave like the free solutions asymptotically as $t \to \pm \infty$, that is, we want to find unique $f_0^{\pm} \in E$ such that

$$||U(t,0)f - U_0(t)f_0^{\pm}||_E \to 0, \quad (t \to \pm \infty)$$

where $U_0(t)$ is the unitary group in E which represents the solution of the free equation (V = 0) and U(t, s) is the evolution operator which maps the solutions at time t to those at time s:

$$u(t) = U(t,s)u(s).$$

In order to state our results, we introduce the following conditions.

(A0) V(x,t) satisfies

$$-\delta \left\{ m^2 + \left(\frac{n-2}{2r}\right)^2 \right\} \le V(x,t) \le C_1 \left\{ m^2 + \left(\frac{n-2}{2r}\right)^2 \right\},$$

where $0 < \delta < 1 \leq C_1$.

(A1) V(x,t) satisfies

$$|V_t(x,t)| \le \eta_1(t) \left\{ m^2 + \left(\frac{n-2}{2r}\right)^2 \right\},\$$

where $\eta_1(t)$ is a positive L^1 -function of $t \ge 0$.

(A2) m(x) and V(x,t) satisfy

$$(m^2\psi)_r \le 0,$$
$$(V\psi)_r \le \{\varepsilon\psi'(r) + \eta_2(t)\} \left(\frac{n-2}{2r}\right)^2,$$

where $\varepsilon > 0$ is sufficiently small, $\eta_2(t)$ is a positive L^1 -function of $t \ge 0$ and $\psi(r)$ is bounded and monotone increasing function of r = |x| such that

$$0 < \psi'(r) < r^{-1}\psi(r), \quad \psi''(r) \le 0.$$

(A3) V(x,t) satisfies

$$|V(x,t)| \le \{C_2\psi'(r) + \eta_3(t)\}\left(\frac{n-2}{2r}\right),\$$

where $\eta_3(t)$ is a positive L^1 -function of $t \ge 0$.

The following lemma plays an important role in our results.

<u>**Lemma 1**</u> (Space-time weighted estimate) Under (A2), there exists $C_3 > 0$ such that for any t > s

$$\int_{s}^{t} \int_{\mathbf{R}^{n}} \left\{ \frac{1}{2} \psi'(r) (u_{\tau}^{2} + |\nabla u|^{2}) - \psi''(r) \frac{n-1}{4r} u^{2} \right\} dx d\tau < C_{3} ||u(s)||_{E}^{2}.$$

By using this lemma, we shall show that energy nondecay and scattering occur.

Theorem 2 (Nondecay)

Assume $(A0) \sim (A3)$,

If $u^{s}(t)$ is the solution of the initial value probrem

$$\begin{cases} \partial_t^2 u^s - \Delta u^s + m^2 u^s + V u^s = 0, \quad t \ge s \\ u^s(s) = u_0(s) \end{cases}$$

and s is sufficiently large, then its energy $||u^s(t)||_E$ never decays as $t \to \infty$.

<u>Theorem 3</u> (Scattering)

Assume $(A0) \sim (A3)$. Then

(i) For every $f \in E$ there exists $f_0^{\pm} \in E$ such that

$$||U(t,s)f - U_0(t-s)f_0^{\pm}||_E^2 \to 0, \quad (t \to \pm \infty)$$

we put

$$Z^{\pm} = s - \lim_{t \to \pm \infty} U_0(-t)U(t,0)$$

Then Z^{\pm} defines a nontrivial bounded operator on E.

(*ii*) If C_2 in (A3) is sufficiently small, then Z^{\pm} gives a bijection on E. Thus, the scattering operator

$$S = Z^+ (Z^-)^{-1} : f_0^- \to f_0^+$$

is well defined and also gives a bijection on E.

§2. Strichartz estimates and its applications

We introduce the Strichartz estimates for the Klein-Gordon equations and as its application, we shall show the unique existence of the solutions of the integral equations corresponded to (P). Define

$$\Phi[f](t) = \int_{s}^{t} e^{iA(t-\tau)} f(\tau) d\tau.$$

Lemma 4 (Strichartz estimates for Klein-Gordon equations) Let $n \ge 3$, $\frac{n-2}{2n} < \frac{1}{p} \le \frac{1}{2}$ and $\frac{1}{r} = \frac{n}{2} \left(\frac{1}{2} - \frac{1}{p}\right)$. Then for any time interval I and for any $s \in \overline{I}$ the following estimates are true

$$\|\Phi[f]\|_{L^{r}(I;L^{p})} \leq C_{4} \|f\|_{L^{r'}(I;H^{2\mu}_{p'})},$$

$$\|\Phi[f]\|_{L^{\infty}(I;L^{2})} \leq C_{4}^{1/2} \|f\|_{L^{r'}(I;H^{\mu}_{p'})},$$

and

$$\|e^{iAt}\phi\|_{L^{r}(I;L^{p})} \leq C_{4}^{1/2}\|\phi\|_{H^{\mu}},$$

where $\frac{1}{r} + \frac{1}{r'} = 1$, $\frac{1}{p} + \frac{1}{p'} = 1$, $\mu = \left(1 + \frac{n}{2}\right)\left(\frac{1}{2} - \frac{1}{p}\right)$ and $A = (m^{2} - \Delta)^{1/2}$.

(B0) V(x,t) satisfies

$$V(x,t) \in L^{\nu}(\mathbf{R};L^q),$$

where

$$0 \le \frac{1}{q} < \frac{2}{n}$$
 and $\frac{1}{\nu} = 1 - \frac{n}{2q}$.

Moreover, V(x, t) satisfies the smallness condition

$$C_4 \|V\|_{L^{\nu}(\mathbf{R};L^q)} < 1$$

where C_4 is a constant given in Lemma 4.

Since this lemma, we obtain the following theorem.

<u>Theorem 5</u> (Unique existence and space-time $L^{p}-L^{q}$ estimates) Let $n \geq 3$ and assume (B0). Then for each $f = {}^{t}(f_{1}, f_{2}) \in E$,

(i) the integral equation

$$u(t) = \dot{K}(t-s)f_1 + K(t-s)f_2 + \int_s^t K(t-\tau)V(\tau)u(\tau)d\tau$$

has a unique solution in $u(t) \in L^r(\mathbf{R}; H_p^{1/2})$. Here $K(t) = A^{-1} \sin(At)$.

(*ii*) This solution belongs to $C(\mathbf{R}; H^1)$ and coincides with U(t, s)f. Moreover, we have

$$\|u\|_{L^{r}(\mathbf{R};H_{p}^{1/2})} \leq \frac{2C_{4}^{1/2}}{1 - C_{4}} \|V\|_{L^{\nu}(\mathbf{R};L^{q})} \|f\|_{E}$$

and

$$\left\| \int_{s}^{t} K(t-\tau)V(\tau)u(\tau)d\tau \right\|_{H^{1}} \leq \frac{2C_{4}\|V\|_{L^{\nu}(\mathbf{R};L^{q})}}{1-C_{4}\|V\|_{L^{\nu}(\mathbf{R};L^{q})}} \|f\|_{E}$$

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